NYX User's Guide

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# Contents

1	Intr	oduction	7
<b>2</b>	Get	ting Started	9
	2.1	Getting Started With Git	9
	2.2	Building the Code	10
	2.3	Running the Code	10
	2.4	Visualization of the Results	10
3	Inp	ıts Files	11
•	3.1	Problem Geometry	11
		3.1.1 List of Parameters	
		3.1.2 Examples of Usage	
	3.2	Domain Boundary Conditions	
	0.2	3.2.1 List of Parameters	
		3.2.2 Notes	
		3.2.3 Examples of Usage	12
	3.3	Resolution	12
	0.0	3.3.1 List of Parameters	12
		3.3.2 Examples of Usage	
	3.4	Tagging	
	0.1	3.4.1 List of Parameters	13
		3.4.2 Notes	13
	3.5	Regridding	13
	0.0	3.5.1 Overview	13
		3.5.2 List of Parameters	
		3.5.3 Notes	
		3.5.4 Examples of Usage	14
		3.5.5 How Grids are Created	15
	3.6	Simulation Time	$15 \\ 15$
	5.0	3.6.1 List of Parameters	
		3.6.2 Notes	
		3.6.3 Examples of Usage	
	3.7	Time Step	
	5.1	1	
		3.7.2 Examples of Usage	10

	3.8	Subcycling
		3.8.1 List of Parameters
		3.8.2 Examples of Usage
	3.9	Restart Capability
		3.9.1 List of Parameters
		3.9.2 Notes
		3.9.3 Examples of Usage
9		Controlling PlotFile Generation
		3.10.1 List of Parameters
		3.10.2 Notes
		3.10.3 Examples of Usage
4		•
•		Screen Output
		3.11.1 List of Parameters
		3.11.2 Notes
		3.11.3 Examples of Usage
٠	3.12	Gravity
		3.12.1 List of Parameters
		3.12.2 Notes
;	3.13	Physics
		3.13.1 List of Parameters
5 ]	Equa	ations in Comoving Coordinates
5 ]	<b>Equ</b> a 5.1	ations in Comoving Coordinates  Hydrodynamic Equations in Comoving Coordinates
5 ]	<b>Equ</b> a 5.1	ations in Comoving Coordinates  Hydrodynamic Equations in Comoving Coordinates
<b>5</b> ]	Equa 5.1	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates
<b>5</b> ]	Equa 5.1	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates
<b>5</b> ]	Equa 5.1	Ations in Comoving Coordinates  Hydrodynamic Equations in Comoving Coordinates
5 ]	Equa 5.1 5.2 Grav	Ations in Comoving Coordinates  Hydrodynamic Equations in Comoving Coordinates
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form  5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations  Equations
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form  5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles  7.2.1 Read from an ASCII file  7.2.2 Read from a binary file
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form  5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates 5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles 7.2.5 Random placement
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates 5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles 7.2.5 Random placement 7.2.6 Cosmological
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form  5.1.2 Tracing  Subgrid Scale Model in Comoving Coordinates  vity  K Matter Particles  Equations  Initializing the Particles  7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary file 7.2.4 Reading SPH particles 7.2.5 Random placement 7.2.6 Cosmological 7.2.6.1 Generating a transfer function
5 ] 6 (7 ]	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates 5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  Vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles 7.2.5 Random placement 7.2.6 Cosmological 7.2.6.1 Generating a transfer function 7.2.6.2 Setting up the initial displacements
5 1 6 6 6 7 1 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates  5.1.1 Conservative Form  5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  Vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles 7.2.5 Random placement 7.2.6 Cosmological 7.2.6.1 Generating a transfer function 7.2.6.2 Setting up the initial displacements 7.2.6.3 Using Nyx with cosmological initial conditions
5 1 5 6 6 6 7 1 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	Equa 5.1 5.2 Grav Dark 7.1 7.2	Ations in Comoving Coordinates Hydrodynamic Equations in Comoving Coordinates 5.1.1 Conservative Form 5.1.2 Tracing Subgrid Scale Model in Comoving Coordinates  Vity  K Matter Particles Equations Initializing the Particles 7.2.1 Read from an ASCII file 7.2.2 Read from a binary file 7.2.3 Read from a binary "meta" file 7.2.4 Reading SPH particles 7.2.5 Random placement 7.2.6 Cosmological 7.2.6.1 Generating a transfer function 7.2.6.2 Setting up the initial displacements

		7.3.2	Motion by	Self-Gra	vity			 	 	 	 			40
			7.3.2.1 N	Move-Kicl	κ-Drift A	Algorith	hm .	 	 	 	 			40
			7.3.2.2	Computin	g <b>g</b>			 	 	 	 			40
	7.4	Outpu	t Format .					 	 	 	 			41
		7.4.1	Checkpoin	nt Files .				 	 	 	 			41
		7.4.2	Plot Files					 	 	 	 			41
		7.4.3	ASCII Par	rticle File	s			 	 	 	 			41
		7.4.4	Run-time	Data Log	S			 	 	 	 			42
		7.4.5	Run-time	Screen O	utput			 	 	 	 			42
8	Visi	ualizati	on											43
	8.1	amrvis						 	 	 	 			43
	8.2	VisIt .						 	 	 	 			43
	8.3													
	8.4		lling What											
9	Soft	ware I	rameworl	k										45
	9.1	Code s	tructure .					 	 	 	 			45
	9.2	Variab	le Names .					 	 	 	 			46
	9.3		l I/O											
10	Veri	ificatio	n Test Pr	oblems										49
	10.1	Cosmo	logy Test I	Problems				 	 	 	 			49
			90Mpc Bo											
Re	efere	nces												49

# Introduction

Welcome to the NYX User's Guide!

In this User's Guide we describe how to download and run Nyx, a massively parallel code that couples the compressible hydrodynamic equations on a grid with a particle representation of dark matter.

Nyx uses an Eulerian grid for the hydrodynamics solver and incorporates adaptive mesh refinement (AMR).

# Getting Started

# 2.1 Getting Started With Git

Please note that Nyx is not yet available for public release. If you have heard about Nyx and are interested in trying it, please contact Ann Almgren at asalmgren@lbl.gov

Nyx is now distributed in two parts using git – you must first download the BoxLib repository, then the Nyx repo.

To use Nyx:

1. Assuming git is installed on your machine – and we recommend version 1.7.x or higher – download the BoxLib repository by typing

git clone https://ccse.lbl.gov/pub/Downloads/BoxLib.git

This will create a directory called BoxLib on your machine. Put this somewhere out of the way and set the environment variable, **BOXLIB\_HOME**, on your machine to the path name where you have put BoxLib. You will want to periodically update BoxLib by typing

git pull

in the BoxLib directory.

2. We will have set up an account for you; follow the instructions we have given you to access Nyx itself.

## 2.2 Building the Code

1. From the directory in which you checked out Nyx, type

cd Nyx/Exec/Test\_90Mpc

This will put you into a directory in which you can run the 90Mpc test problem from http://t8web.lanl.gov/people/heitmann/arxiv/

2. In Test\_90Mpc, edit the GNUmakefile, and set

COMP = your favorite C++ compiler

FCOMP = your favorite Fortran compiler (which must compile F90)

DEBUG = FALSE

We like COMP = gcc and FCOMP = gfortran.

3. Now type "make". The resulting executable will look something like "Nyx3d.Linux.gcc.gfortran.ex", which means this is a 3-d version of the code, made on a Linux machine, with COMP = gcc and FCOMP = gfortran.

## 2.3 Running the Code

- 1. Type "Nyx3d.Linux.gcc.gfortran.ex inputs"
- 2. You will notice that running the code generates directories that look like plt00000, plt00020, etc, and chk00000, chk00020, etc. These are "plotfiles" and "checkpoint" files. The plotfiles are used for visualization, the checkpoint files are used for restarting the code.

## 2.4 Visualization of the Results

1. To visualize the plotfiles, you can use Visit, yt, or a homegrown visualization tool known as amrvis.

To use **amrvis**, please contact Ann Almgren (asalmgren@lbl.gov) for information on how to download it. Visit and yt are available from their web sites.

You have now completed a brief introduction to Nyx.

# Inputs Files

The Nyx executable reads run-time information from an "inputs" file (which you put on the command line) and from a "probin" file, the name of which is usually defined in the inputs file, but which defaults to "probin". To set the "probin" file name in the inputs file:

amr.probin\_file = my\_special\_probin

for example, has the Fortran code read a file called "my\_special\_probin"

# 3.1 Problem Geometry

## 3.1.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
geometry.prob_lo	physical location of low corner of the domain	Real	must be set
geometry.prob_hi	physical location of high corner of the domain	Real	must be set
geometry.coord_sys	coordinate system	0 = Cartesian, 1 = r-z, 2 = spherical	must be set
geometry.is_periodic	is the domain periodic in this direction	0 if false, 1 if true	0 0 0

## 3.1.2 Examples of Usage

- **geometry.prob\_lo** = 0 0 0 defines the low corner of the domain at (0,0,0) in physical space.
- **geometry.prob\_hi** = 1.e8 2.e8 2.e8 defines the high corner of the domain at (1.e8,2.e8,2.e8) in physical space.
- **geometry.coord\_sys** = 0 defines the coordinate system as Cartesian
- **geometry.is\_periodic** = 0 1 0 says the domain is periodic in the y-direction only.

# 3.2 Domain Boundary Conditions

## 3.2.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
nyx.lo_bc	boundary type of each low face	0,1,2,3,4,5	must be set
nyx.hi_bc	boundary type of each high face	0,1,2,3,4,5	must be set

### 3.2.2 Notes

Boundary types are:

0 – Interior / Periodic	3 – Symmetry
1 - Inflow	4 – Slip Wall
2 – Outflow	5 – No Slip Wall

Note – nyx.lo\_bc and nyx.hi\_bc must be consistent with **geometry.is\_periodic** – if the domain is periodic in a particular direction then the low and high bc's must be set to 0 for that direction.

## 3.2.3 Examples of Usage

- $nyx.lo_bc = 140$
- $nyx.hi_bc = 240$
- geometry.is\_periodic = 0 0 1

would define a problem with inflow (1) in the low-x direction, outflow(2) in the high-x direction, slip wall (4) on the low and high y-faces, and periodic in the z-direction.

## 3.3 Resolution

## 3.3.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.n_cell	number of cells in each direction at the coarsest level	Integer $> 0$	must be set
amr.max_level	number of levels of refinement above the coarsest level	Integer $\geq 0$	must be set
amr.ref_ratio	ratio of coarse to fine grid spacing between subsequent levels	2 or 4	must be set
amr.regrid_int	how often to regrid	Integer $> 0$	must be set
amr.regrid_on_restart	should we regrid immediately after restarting	0 or 1	0

Note: if amr.max\_level = 0 then you do not need to set amr.ref\_ratio or amr.regrid\_int.

## 3.3.2 Examples of Usage

### • $amr.n_cell = 32 64 64$

would define the domain to have 32 cells in the x-direction, 64 cells in the y-direction, and 64 cells in the z-direction at the coarsest level. (If this line appears in a 2D inputs file then the final number will be ignored.)

#### • $amr.max\_level = 2$

would allow a maximum of 2 refined levels in addition to the coarse level. Note that these additional levels will only be created only if the tagging criteria are such that cells are flagged as needing refinement. The number of refined levels in a calculation must be  $\leq$  amr.max\_level, but can change in time and need not always be equal to amr.max\_level.

## • $amr.ref_ratio = 24$

would set factor 2 refinement between levels 0 and 1, and factor 4 refinement between levels 1 and 2. Note that you must have at least **amr.max\_level** values of **amr.ref\_ratio** (Additional values may appear in that line and they will be ignored).

### • amr.regrid\_int = $2\ 2$

tells the code to regrid every 2 steps. Thus in this example, new level 1 grids will be created every 2 level 0 time steps, and new level 2 grids will be created every 2 level 1 time steps.

## 3.4 Tagging

## 3.4.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
nyx.allow_untagging	are cells allowed to be "untagged"	0 or 1	0

### 3.4.2 Notes

• Typically cells at a given level can be tagged as needing refinement by any of a number of criteria, but cannot be "untagged", i.e. once tagged no other criteria can untag them. If we set nyx.allow\_untagging = 1 then the user is allowed to "untag" cells in the Fortran tagging routines

# 3.5 Regridding

#### 3.5.1 Overview

The details of the regridding strategy are described in a later section; here we cover how the input parameters can control the gridding.

As described later, the user defines Fortran subroutines which tag individual cells at a given level if they need refinement. This list of tagged cells is sent to a grid generation routine, which uses the Berger-Rigoutsis algorithm to create rectangular grids that contain the tagged cells.

### 3.5.2 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.regrid_file	name of file from which to read the grids	text	no file
amr.grid_eff	grid efficiency at coarse level at which grids are created	Real > 0 and $< 1$	0.7
amr.n_error_buf	radius of additional tagging around already tagged cells	Integer $\geq 0$	1
amr.max_grid_size	maximum size of a grid in any direction	Integer $> 0$	128 in 2D, 32 in 3D
amr.blocking_factor	grid size must be a multiple of this	Integer $> 0$	2
amr.refine_grid_layout	refine grids more if # of processors > # of grids	0 if false, 1 if true	1

#### 3.5.3 Notes

- amr.n\_error\_buf, amr.max\_grid\_size and amr.blocking\_factor can be read in as a single value which is assigned to every level, or as multiple values, one for each level
- amr.max\_grid\_size at every level must be even
- amr.blocking\_factor at every level must be a power of 2
- the domain size must be a multiple of amr.blocking\_factor at level 0
- amr.max\_grid\_size must be a multiple of amr.blocking\_factor at every level

## 3.5.4 Examples of Usage

### • amr.regrid\_file = fixed\_grids

In this case the list of grids at each fine level are contained in the file, fixed\_grids, which will be read during the gridding procedure. These grids must not violate the amr.max\_grid\_size criterion. The rest of the gridding procedure described below will not occur if amr.regrid\_file is set.

## • $\operatorname{amr.grid\_eff} = 0.9$

During the grid creation process, at least 90% of the cells in each grid at the level at which the grid creation occurs must be tagged cells. Note that this is applied at the coarsened level at which the grids are actually made, and before **amr.max\_grid\_size** is imposed.

## • $amr.max\_grid\_size = 64$

The final grids will be no longer than 64 cells on a side at every level.

#### • amr.max\_grid\_size = 64 32 16

The final grids will be no longer than 64 cells on a side at level 0, 32 cells on a side at level 1, and 16 cells on a side at level 2.

### • amr.blocking\_factor = 32

The dimensions of all the final grids will be multiples of 32 at all levels.

#### • amr.blocking\_factor = 32 16 8

The dimensions of all the final grids will be multiples of 32 at level 0, multiples of 16 at level 1, and multiples of 8 at level 2..

Having grids that are large enough to coarsen multiple levels in a V-cycle is essential for good multigrid performance in simulations that use self-gravity.

### 3.5.5 How Grids are Created

The gridding algorithm proceeds in this order:

- 1. Grids are created using the Berger-Rigoutsis clustering algorithm modified to ensure that all new fine grids are divisible by **amr.blocking\_factor**.
- 2. Next, the grid list is chopped up if any grids are larger than max\_grid\_size. Note that because amr.max\_grid\_size is a multiple of amr.blocking\_factor the amr.blocking\_factor criterion is still satisfied.
- 3. Next, if  $amr.refine\_grid\_layout = 1$  and there are more processors than grids, and
  - if amr.max\_grid\_size / 2 is a multiple of amr.blocking\_factor

then the grids will be redefined, at each level independently, so that the maximum length of a grid at level  $\ell$ , in any dimension, is **amr.max\_grid\_size**[ $\ell$ ] / 2.

- 4. Finally, if amr.refine\_grid\_layout = 1, and there are still more processors than grids, and
  - if amr.max\_grid\_size / 4 is a multiple of amr.blocking\_factor

then the grids will be redefined, at each level independently, so that the maximum length of a grid at level  $\ell$ , in any dimension, is **amr.max\_grid\_size**[ $\ell$ ] / 4.

## 3.6 Simulation Time

## 3.6.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
max_step	maximum number of level 0 time steps	Integer $\geq 0$	-1
$stop\_time$	final simulation time	$Real \ge 0$	-1.0
nyx.final_a	if $nyx.use\_comoving = t$ and positive value then this is final value of a	Real > 0	-1.0
nyx.final_z	if $nyx.use\_comoving = t$ and positive value then this is final value of z	Real > 0	-1.0

## 3.6.2 Notes

To control the number of time steps, you can limit by the maximum number of level 0 time steps (max\_step), or the final simulation time (stop\_time), or both. The code will stop at whichever criterion comes first. Note that if the code reaches stop\_time then the final time step will be shortened so as to end exactly at stop\_time, not pass it.

If running in comoving coordinates you can also set a final value of a by setting **nyx.final\_a**, or a final value of z by setting **nyx.final\_z**. You may only specify one or the other of these. Once this value of a or z is reached in a time step, the code will stop at the end of this full coarse time step. (Note it does not stop at a (or z) exactly equal to the final value, rather it stops once a is greater than (or z is less than) this value.)

## 3.6.3 Examples of Usage

- $max\_step = 1000$
- $stop\_time = 1.0$

will end the calculation when either the simulation time reaches 1.0 or the number of level 0 steps taken equals 1000, whichever comes first.

# 3.7 Time Step

- If  $\mathbf{nyx.do\_hydro} = 1$ , then typically the code chooses a time step based on the CFL number (dt = cfl \* dx / max(u+c)).
- If  $\mathbf{nyx.do\_hydro} = 0$  and we are running with dark matter particles, then we use a time step based on the velocity of the particles, i.e. we set  $\Delta t$  so that the particle goes no further than  $f\Delta t$  in a coordinate direction where  $0 \le f \le 1$ . The value for f is currently hard-wired in Particles.H, but it will become an inputs parameter.

### 3.7.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
nyx.cfl	CFL number for hydro	Real $> 0$ and $\le 1$	0.8
particles.cfl	CFL number for particles	Real $> 0$ and $\le 1$	0.5
nyx.init_shrink	factor by which to shrink the initial time step	Real $> 0$ and $\le 1$	1.0
nyx.change_max	factor by which the time step can grow in subsequent steps	$Real \ge 1$	1.1
$nyx.fixed_dt$	level 0 time step regardless of cfl or other settings	Real > 0	unused if not set
nyx.initial_dt	initial level 0 time step regardless of other settings	Real > 0	unused if not set
$nyx.dt_cutoff$	time step below which calculation will abort	Real > 0	0.0

## 3.7.2 Examples of Usage

- nyx.cfl = 0.9 defines the timestep as dt = cfl \* dx / umax\_hydro.
- particles.cfl = 0.9 defines the timestep as dt = cfl \* dx / umax\_particles where umax\_particles is the maximum velocity of any particle in the domain.
- nyx.init\_shrink = 0.01 sets the initial time step to 1% of what it would be otherwise.
- nyx.change\_max = 1.1 allows the time step to increase by no more than 10% in this case. Note that the time step can shrink by any factor; this only controls the extent to which it can grow.
- nyx.fixed\_dt = 1.e-4
  sets the level 0 time step to be 1.e-4 for the entire simulation, ignoring the other timestep controls. Note that if nyx.init\_shrink ≠ 1 then the first time step will in fact be nyx.init\_shrink \* nyx.fixed\_dt.

## • $nyx.initial_dt = 1.e-4$

sets the *initial* level 0 time step to be 1.e-4 regardless of **nyx.cfl** or **nyx.fixed\_dt**. The time step can grow in subsequent steps by a factor of **nyx.change\_max** each step.

### • $\mathbf{nyx.dt\_cutoff} = 1.e-20$

tells the code to abort if the time step ever gets below 1.e-20. This is a safety mechanism so that if things go nuts you don't burn through your entire computer allocation because you don't realize the code is misbehaving.

## 3.8 Subcycling

Nyx supports a number of different modes for subcycling in time.

- If amr.subcycling\_mode=Auto (default), then the code will run with equal refinement in space and time. In other words, if level n + 1 is a factor of 2 refinement above level n, then n + 1 will take 2 steps of half the duration for every level n step.
- If  $amr.subcycling\_mode=$ None, then the code will not refine in time. All levels will advance together with a timestep dictated by the level with the strictest dt. Note that this is identical to the deprecated command amr.nosub = 1.
- If amr.subcycling\_mode=Manual, then the code will subcycle according to the values supplied by subcycling\_iterations.

## 3.8.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.subcycling_mode	How shall we subcycle	Auto, None or Manual	Auto
amr.subcycling_iterations	Number of cycles at each level	1 or ref_ratio	must be set in Manual mode

## 3.8.2 Examples of Usage

## $\bullet \ \mathbf{amr.subcycling\_mode} {=} \mathbf{Manual}$

Subcycle in manual mode with largest allowable timestep.

### • amr.subcycling\_iterations = 1 2 1 2

Take 1 level 0 timestep at a time (required). Take 2 level 1 timesteps for each level 0 step, 1 timestep at level 2 for each level 1 step, and take 2 timesteps at level 3 for each level 2 step.

## • amr.subcycling\_iterations = 2

Alternative form. Subcycle twice at every level (except level 0).

# 3.9 Restart Capability

Nyx has a standard sort of checkpointing and restarting capability. In the inputs file, the following options control the generation of checkpoint files (which are really directories):

## 3.9.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.check_file	prefix for restart files		"chk"
amr.check_int	ck_int how often (by level 0 time steps) to write restart files		-1
amr.check_per	amr.check_per how often (by simulation time) to write restart files		-1.0
amr.restart	amr.restart name of the file (directory) from which to restart		not used if not set
amr.checkpoint_files_output	mr.checkpoint_files_output   should we write checkpoint files		1
amr.check_nfiles	how parallel is the writing of the checkpoint files		64
amr.checkpoint_on_restart should we write a checkpoint immediately after restarting		0 or 1	0

### 3.9.2 Notes

- You should specify either amr.check\_int or amr.check\_per. Do not try to specify both.
- Note that if **amr.check\_per** is used then in order to hit that exact time the code may modify the time step slightly, which will change your results ever so slightly than if you didn't set this flag.
- Note that amr.plotfile\_on\_restart and amr.checkpoint\_on\_restart only take effect if amr.regrid\_on\_restart is in effect.
- See the Software Section for more details on parallel I/O and the **amr.check\_nfiles** parameter.
- If you are doing a scaling study then set **amr.checkpoint\_files\_output** = 0 so you can test scaling of the algorithm without I/O.

## 3.9.3 Examples of Usage

- amr.check\_file = chk\_run
- $amr.check_int = 10$

means that restart files (really directories) starting with the prefix "chk\_run" will be generated every 10 level 0 time steps. The directory names will be  $chk_run000000$ ,  $chk_run00010$ ,  $chk_run00020$ , etc.

If instead you specify

- amr.check\_file = chk\_run
- $amr.check_per = 0.5$

then restart files (really directories) starting with the prefix "chk\_run" will be generated every 0.1 units of simulation time. The directory names will be  $chk_run00000$ ,  $chk_run00043$ ,  $chk_run00061$ , etc, where t = 0.1 after 43 level 0 steps, t = 0.2 after 61 level 0 steps, etc.

To restart from *chk\_run00061*, for example, then set

•  $amr.restart = chk_run00061$ 

# 3.10 Controlling PlotFile Generation

The main output from Nyx is in the form of plotfiles (which are really directories). The following options in the inputs file control the generation of plotfiles

### 3.10.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.plot_file	prefix for plotfiles	Text	"plt"
$amr.plot\_int$	how often (by level 0 time steps) to write plot files	Integer > 0	-1
amr.plot_per	how often (by simulation time) to write plot files	Real > 0	-1.0
$amr.plot\_vars$	name of state variables to include in plotfiles	ALL, NONE or list	ALL
amr.derive_plot_vars	derive_plot_vars name of derived variables to include in plotfiles		NONE
amr.plot_files_output	r.plot_files_output should we write plot files		1
amr.plotfile_on_restart	r.plotfile_on_restart should we write a plotfile immediately after restarting		0
amr.plot_nfiles	plot_nfiles how parallel is the writing of the plotfiles		64
nyx.plot_phiGrav	ot_phiGrav Should we plot the gravitational potential		0
	plot the gravitational potential	0 or 1	0
particles.write_in_plotfile	es.write_in_plotfile   Should we write the particles in a file within the plotfile		0
fab.format	Should we write the plotfile in double or single precision	NATIVE or IEEE32	NATIVE

All the options for **amr.derive\_plot\_vars** are kept in **derive\_lst** in Nyx\_setup.cpp. Feel free to look at it and see what's there.

### 3.10.2 Notes

- You should specify either amr.plot\_int or amr.plot\_per. Do not try to specify both.
- Note that if **amr.plot\_per** is used then in order to hit that exact time the code may modify the time step slightly, which will change your results ever so slightly than if you didn't set this flag.
- See the Software Section for more details on parallel I/O and the amr.plot\_nfiles parameter.
- If you are doing a scaling study then set **amr.plot\_files\_output** = 0 so you can test scaling of the algorithm without I/O.
- nyx.plot\_phiGrav is only relevant if nyx.do\_grav = 1 and gravity.gravity\_type = PoissonGrav
- By default, plotfiles are written in double precision (NATIVE format). If you want to save space by writing them in single precision, set the fab.format flag to IEEE32.

## 3.10.3 Examples of Usage

- $amr.plot_file = plt_run$
- $\operatorname{amr.plot\_int} = 10$

means that plot files (really directories) starting with the prefix "plt\_run" will be generated every 10 level 0 time steps. The directory names will be plt\_run00000, plt\_run00010, plt\_run00020, etc.

If instead you specify

- amr.plot\_file = plt\_run
- $amr.plot_per = 0.5$

then restart files (really directories) starting with the prefix "plt\_run" will be generated every 0.1 units of simulation time. The directory names will be  $plt_run00000$ ,  $plt_run00043$ ,  $plt_run00061$ , etc, where t = 0.1 after 43 level 0 steps, t = 0.2 after 61 level 0 steps, etc.

# 3.11 Screen Output

## 3.11.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
amr.v	verbosity of Amr.cpp	0 or 1	0
nyx.v	verbosity of Nyx.cpp	0 or 1	0
gravity.v	verbosity of Gravity.cpp	0 or 1	0
mg.v	verbosity of multigrid solver (for gravity)	0,1,2,3,4	0
particles.v	verbosity of particle-related processes	0,1,2,3,4	0
amr.grid_log	name of the file to which the grids are written	Text	not used if not set
amr.run_log	name of the file to which certain output is written	Text	not used if not set
amr.run_log_terse	name of the file to which certain (terser) output is written	Text	not used if not set
amr.sum_interval	if $> 0$ , how often (in level 0 time steps)		
	to compute and print integral quantities	Integer	-1
nyx.do_special_tagging		0 or 1	1

#### 3.11.2 Notes

• nyx.do\_special\_tagging = 1 allows the user to set a special flag based on user-specified criteria. This can be used, for example, to calculate the bounce time in a core collapse simulation; the bounce time is defined as the first time at which the maximum density in the domain exceeds a user-specified value. This time can then be printed into a special file as a useful diagnostic.

## 3.11.3 Examples of Usage

•  $\operatorname{\mathbf{amr.grid\_log}} = \operatorname{grdlog}$ 

Every time the code regrids it prints a list of grids at all relevant levels. Here the code will write these grids lists into the file *qrdloq*.

•  $\operatorname{\mathbf{amr.run\_log}} = \operatorname{runlog}$ 

Every time step the code prints certain statements to the screen (if  $\mathbf{amr.v} = 1$ ), such as STEP = 1 TIME = 1.91717746 DT = 1.91717746

PLOTFILE: file = plt00001

Here these statements will be written into runlog as well.

 $\bullet \ \mathbf{amr.run\_log\_terse} = \mathrm{runlogterse}$ 

This file, runlogterse differs from runlog, in that it only contains lines of the form

10 0.2 0.005

in which "10" is the number of steps taken, "0.2" is the simulation time, and "0.005" is the level 0 time step. This file can be plotted very easily to monitor the time step.

## • $nyx.sum_interval = 2$

if nyx.sum\_interval > 0 then the code computes and prints certain integral quantities, such as total mass, momentum and energy in the domain every nyx.sum\_interval level 0 steps. In this example the code will print these quantities every two coarse time steps. The print statements have the form

TIME = 1.91717746 MASS = 1.792410279e + 34

for example. If this line is commented out then it will not compute and print these quantitities.

## 3.12 Gravity

## 3.12.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
nyx.do_grav	Include gravity as a forcing term	0 if false, 1 if true	must be set
gravity.gravity_type	if $\mathbf{nyx.do\_grav} = 1$ , how shall gravity be calculated	CompositeGrav,	
		PoissonGrav	must be set
gravity.no_sync	if gravity.gravity_type = PoissonGrav, whether to perform the "sync solve"	0 or 1	0
gravity.no_composite	if <b>gravity_type</b> = PoissonGrav, whether to perform a composite solve	0 or 1	0

### 3.12.2 Notes

Gravity types are CompositeGrav or PoissonGrav. See the Gravity section for more detail.

- To include gravity you must set
  - USE\_GRAV = TRUE in the GNUmakefile
  - **nyx.do**\_**grav** = 1 in the inputs file
- gravity\_gravity\_type is only relevant if nyx.do\_grav = 1
- gravity.no\_sync and gravity.no\_composite are only relevant if gravity.gravity\_type = PoissonGrav, i.e. the code does a full Poisson solve for self-gravity.

# 3.13 Physics

## 3.13.1 List of Parameters

Parameter	Definition	Acceptable Values	Default
nyx.do_hydro	Time-advance the fluid dynamical equations	0 if false, 1 if true	must be set
nyx.do_react	react Include reactions		must be set
$nyx.add_ext_src$	nyx.add_ext_src Include additional user-specified source term		0
nyx.use_const_species	nst_species If 1 then read h_species and he_species		0
nyx.h_species	Concentration of H	0 < X < 1	0
nyx.he_species	Concentration of He	0 < X < 1	0

# Units and Constants

## 4.1 Units and Constants

We support two different systems of units in Nyx: CGS and Cosmological. All inputs and problem initialization should be specified consistently with one of these sets of units. No internal conversions of units occur within the code, so the output must be interpreted appropriately.

The default is cosmological units.

If you want to use CGS units instead, then set

 $USE\_CGS = TRUE$ 

in your GNUmakefile. This will select the file constants\_cgs.f90 instead of constants\_cosmo.f90 from the Nyx/constants directory.

Location	Variable	CGS	Cosmological	Conversion Data
inputs file	geometry.prob_lo geometry.prob_hi	cm cm	Mpc Mpc	$1 \mathrm{Mpc} = 3.08568025e24 \ \mathrm{cm}$ $1 \mathrm{Mpc} = 3.08568025e24 \ \mathrm{cm}$
Hydro Initialization	density	$\rm g \ / \ cm^3$	${ m M}_{\odot}$ / ${ m Mpc}^3$	$1 (M_{\odot} / Mpc^3) = .06769624e-39 (g/cm^3)$
Hydro Initialization	velocities	cm/s	km/s	$1 \mathrm{km} = 1.\mathrm{e}5 \mathrm{~cm}$
Hydro Initialization	momenta	(g/cm <sup>3</sup> ) (cm/s)	$({\rm M}_{\odot}/{\rm Mpc^3})~({\rm km/s})$	$1 \mathrm{km} = 1.e5 \ \mathrm{cm}$ $1 \ (\mathrm{M}_{\odot} \ / \ \mathrm{Mpc}^3) = .06769624 \mathrm{e}\text{-}39 \ \mathrm{g/cm}^3$
Hydro Initialization	temperature	К	К	1
Hydro Initialization	specific energy $(e \text{ or } E)$	$erg/g=(cm/s)^2$	$(\mathrm{km/s})^2$	$1 (km/s)^2 = 1.e10 (cm/s)^2$
Hydro Initialization	energy $(\rho e \text{ or } \rho E)$	$erg / cm^3 = (g/cm^3) (cm/s)^2$	$({\rm M}_{\odot}/{\rm Mpc^3})~({\rm km/s})^2$	$\begin{array}{c} 1~(km/s)^2 = 1.e10~(cm/s)^2 \\ 1~(M_{\odot}~/~Mpc^3) = .06769624e\text{-}39~g/cm^3 \end{array}$
Particle Initialization	particle mass	g	${ m M}_{\odot}$	$1 {\rm ~M}_{\odot} = 1.98892e33 {\rm ~g}$
Particle Initialization	particle locations	cm	Мрс	1  Mpc = 3.08568025e24  cm
Particle Initialization	particle velocities	cm/s	km/s	1  km = 1e5  cm
Output	Pressure	$g (cm/s)^2 / cm^3$	${\rm M}_{\odot}~({\rm km/s})^2~/~{\rm Mpc}^3$	$1 \text{ M}_{\odot} \text{ (km/s)}^2 / \text{Mpc}^3 = .06769624\text{e-}29 \text{ g (cm/s)}^2 / \text{cm}^3$
Output	Gravity	(cm/s) / s	$(\mathrm{km/s})^2$ / Mpc	$1~{\rm M}_{\odot}~({\rm km/s})^2~/~{\rm Mpc}^3 =$
Output	Time	s	(Mpc/km) s	1  Mpc = 3.08568025e19  km

Table 4.1: Units

Constant	CGS	Cosmological	Conversion Data
Gravitational constant (G)	$6.67428e-8 \text{ cm } (\text{cm/s})^2 \text{ g}^{-1}$	$4.3019425$ e-9 Mpc $(\text{km/s})^2 \text{ M}_{\odot}^{-1}$	
Avogadro's number $(n_A)$	$6.02214129e23 \text{ g}^{-1}$	$1.1977558e57~{ m M}_{\odot}^{-1}$	$1 {\rm ~M}_{\odot} = 1.98892e33 {\rm ~g}$
Boltzmann's constant $(k_B)$	1.3806488e-16 erg / K	$0.6941701$ e-59 M $_{\odot}~({\rm km/s})^2~/~{\rm K}$	$1 \text{ M}_{\odot} \text{ (km/s)}^2 = 1.98892e43 g (cm/s)^2$
Hubble constant $(H)$	100 (km/s) / Mpc	$32.407764868e-19 s^{-1}$	1  Mpc = 3.08568025e19  km

Table 4.2: Constants

The only other place that dimensional numbers are used in the code is in the tracing and Riemann solve. We set three *small* numbers which need to be consistent with the data specified. Each of these can be specified in the inputs file.

- small\_dens small value for density
- small\_p small value for pressure
- small\_T small value for temperature

These are the places that each is used in the code:

### • small\_dens

- subroutine enforce\_minimum\_density (called after subroutine consup) – if  $\rho$  < small\_dens then  $\rho$  is set to the minimum value of the 26 neighbors. This also modifies momenta,  $\rho E$  and  $\rho e$  so that velocties, E and e remain unchanged.

```
    subroutine tracexy / tracez / tracexy_ppm / tracez_ppm:
    qxp = max(qxp,small_dens)
    qxm = max(qxm,small_dens)
    and analogously for qyp/qym and qzp/qzm. This only modifies density inside the tracing, not the other variables
```

- subroutine riemannus - we set

```
wsmall = small_dens * csmall and then wl = max(wsmall, sqrt(gaml * pl * rl)) wr = max(wsmall, sqrt(gamr * pr * rr)) Also, we set ro = max(small_dens,ro) where ro = 0.5 * (rl + rr) - this state is only chosen when ustar = 0, and rstar = max(small_dens,rstar) where rstar = ro + (pstar-po)/co<sup>2</sup>
```

- subroutine react\_state - only compute reaction if  $\rho > \text{small\_dens}$ 

### • small\_temp:

- subroutine ctoprim: if  $\rho e < 0$ , then

call subroutine nyx\_eos\_given\_RTX (e,...,small\_temp,...) in order to compute a new energy, e.

This energy is then used to

call subroutine nyx\_eos\_given\_ReX in order to compute the sound speed, c.

Coming out of this the temperature is equal to small\_temp and the energy e has been reset.

- subroutine react\_state: if  $\rho e < 0$ , then

call subroutine nyx\_eos\_given\_RTX (e,...,small\_temp,...) in order to compute a new energy, e.

This energy is then used to proceed with the burner routine.

- subroutine reset\_internal\_energy: if e < 0 and E - ke < 0 then

call subroutine nyx\_eos\_given\_RTX (e,...,small\_temp,...) in order to compute a new energy, e. This energy is also used to

define a new E = e + ke

## • small\_pres:

- subroutine riemannus - we set

 $pstar = max(small\_pres,pstar)$ 

pgdnv = max(small\_pres,pgdnv). Note that pgdnv is the pressure explicitly used in the fluxes.

- subroutine uflaten small\_pres is used to keep the denominator away from zero
- Everywhere we define values of pressure on a face, we set that value to be at least small\_pres.

# **Equations in Comoving Coordinates**

## 5.1 Hydrodynamic Equations in Comoving Coordinates

## 5.1.1 Conservative Form

We solve the equations of gas dynamics in a coordinate system that is comoving with the expanding universe, with expansion factor, a, related to the redshift, z, by a = 1/(1+z). The continuity equation is written,

$$\frac{\partial \rho_b}{\partial t} = -\frac{1}{a} \nabla \cdot (\rho_b \mathbf{U}) , \qquad (5.1)$$

where  $\rho_b$  is the comoving baryonic density, related to the proper density by  $\rho_b = a^3 \rho_{proper}$ , and U is the proper peculiar baryonic velocity.

The momentum evolution equation can be expressed as

$$\frac{\partial(\rho_b \mathbf{U})}{\partial t} = \frac{1}{a} \left( -\nabla \cdot (\rho_b \mathbf{U} \mathbf{U}) - \nabla p + \rho_b \mathbf{g} + \mathbf{S}_{\rho \mathbf{U}} - \dot{a} \rho_b \mathbf{U} \right) , \qquad (5.2)$$

or equivalently,

$$\frac{\partial(a\rho_b\mathbf{U})}{\partial t} = -\nabla \cdot (\rho_b\mathbf{U}\mathbf{U}) - \nabla p + \rho_b\mathbf{g} + \mathbf{S}_{\rho\mathbf{U}} , \qquad (5.3)$$

where the pressure, p, that appears in the evolution equations is related to the proper pressure,  $p_{proper}$ , by  $p = a^3 p_{proper}$ . Here  $\mathbf{g} = -\nabla \phi$  is the gravitational acceleration vector, and  $\mathbf{S}_{\rho \mathbf{U}}$  represents any external forcing terms.

The energy equation can be written,

$$\frac{\partial(\rho_b E)}{\partial t} = \frac{1}{a} \left[ -\nabla \cdot (\rho_b \mathbf{U} E + p \mathbf{U}) + (\rho_b \mathbf{U} \cdot \mathbf{g} + S_{\rho E}) - \dot{a} (3(\gamma - 1)\rho_b e + \rho_b (\mathbf{U} \cdot \mathbf{U})) \right] . \tag{5.4}$$

or equivalently,

$$\frac{\partial (a^2 \rho_b E)}{\partial t} = a \left[ -\nabla \cdot (\rho_b \mathbf{U} E + p \mathbf{U}) + \rho_b \mathbf{U} \cdot \mathbf{g} + S_{\rho E} + \dot{a} ((2 - 3(\gamma - 1)) \rho_b e) \right] . \tag{5.5}$$

Here  $E=e+\mathbf{U}\cdot\mathbf{U}/2$  is the total energy per unit mass, where e is the specific internal energy.  $S_{\rho E}=S_{\rho e}+\mathbf{U}\cdot\mathbf{S}_{\rho \mathbf{U}}$  where  $S_{\rho e}=\Lambda^H-\Lambda^C$  represents the heating and cooling terms, respectively. We can write the evolution equation for internal energy as

$$\frac{\partial(\rho_b e)}{\partial t} = \frac{1}{a} \left[ -\nabla \cdot (\rho_b \mathbf{U} e) - p\nabla \cdot \mathbf{U} - \dot{a}(3(\gamma - 1)\rho_b e) + S_{\rho e} \right] . \tag{5.6}$$

or equivalently,

$$\frac{\partial (a^2 \rho_b e)}{\partial t} = a \left[ -\nabla \cdot (\rho_b \mathbf{U} e) - p \nabla \cdot \mathbf{U} + S_{\rho e} + \dot{a} ((2 - 3(\gamma - 1)) \rho_b e) \right] . \tag{5.7}$$

Note that for a gamma-law gas with  $\gamma = 5/3$ , we can write

$$\frac{\partial (a^2 \rho_b E)}{\partial t} = a \left[ -\nabla \cdot (\rho_b \mathbf{U} E + p \mathbf{U}) + \rho_b \mathbf{U} \cdot \mathbf{g} + S_{\rho e} \right] . \tag{5.8}$$

and

$$\frac{\partial (a^2 \rho_b e)}{\partial t} = a \left[ -\nabla \cdot (\rho_b \mathbf{U} e) - p \nabla \cdot \mathbf{U} + S_{\rho e} \right] . \tag{5.9}$$

## 5.1.2 Tracing

In order to compute the fluxes on faces, we trace  $\rho$ ,  $\mathbf{U}$ ,  $\rho e$  and p to the faces.

Thus we must convert the momentum evolution equation into a velocity evolution equation:

$$\frac{\partial \mathbf{U}}{\partial t} = \frac{1}{\rho_b} \left( \frac{\partial (\rho_b \mathbf{U})}{\partial t} - \mathbf{U} \frac{\partial \rho_b}{\partial t} \right)$$
 (5.10)

$$= \frac{1}{a\rho_b} \left( -\nabla \cdot (\rho_b \mathbf{U} \mathbf{U}) - \nabla p + \rho_b \mathbf{g} + S_{\rho \mathbf{U}} - \dot{a}\rho_b \mathbf{U} \right) + \frac{1}{a} \mathbf{U} \nabla \cdot (\rho_b \mathbf{U})$$
 (5.11)

$$= \frac{1}{a} \left( -\mathbf{U} \cdot \nabla \mathbf{U} - \frac{1}{\rho_b} \nabla p + \mathbf{g} + \frac{1}{\rho_b} \mathbf{S}_{\rho \mathbf{U}} - \dot{a} \mathbf{U} \right) . \tag{5.12}$$

# 5.2 Subgrid Scale Model in Comoving Coordinates

The fundamental modification to the standard compressible equations is the addition of a SGS turbulence energy variable, K and associated source terms in the equations for the evolution of velocity, total energy, and K [5, 2, 4]. The set of conservation equations in comoving coordinates (5.1)–(5.5) then becomes [1]:

$$\frac{\partial \rho_b}{\partial t} = -\frac{1}{a} \nabla \cdot (\rho_b \mathbf{U}) , \qquad (5.13)$$

$$\frac{\partial(a\rho_b\mathbf{U})}{\partial t} = -\nabla \cdot (\rho_b\mathbf{U}\mathbf{U}) - \nabla p + \nabla \cdot \boldsymbol{\tau} + \rho_b\mathbf{g} , \qquad (5.14)$$

$$\frac{\partial (a^2 \rho_b E)}{\partial t} = -a \nabla \cdot (\rho_b \mathbf{U} E + p \mathbf{U}) + a \rho_b \mathbf{U} \cdot \mathbf{g} + a \nabla \cdot (\mathbf{U} \cdot \boldsymbol{\tau}) - a^2 (\Sigma - \rho_b \varepsilon)$$

$$+ a \dot{a} \left( (2 - 3(\gamma - 1)) \rho_b e \right) + a^2 (\Lambda^H - \Lambda^C) ,$$
(5.15)

$$\frac{\partial (a^2 \rho_b K)}{\partial t} = -a \nabla \cdot (\rho_b \mathbf{U} K) + a \nabla \cdot (\rho_b \kappa_{\text{sgs}} \nabla K) + a^2 (\Sigma - \rho_b \varepsilon) . \tag{5.16}$$

The interaction between resolved and unresolved turbulent eddies is described by the SGS turbulence stress tensor  $\tau$ . Since inertial-range dynamics of turbulence is scale-invariant, we conjecture that  $\tau$  in comoving coordinates has the same form as for non-expanding fluids. For compressible turbulence, the following closure is proposed in [3]:

$$\tau_{ij} = 2C_1 \Delta \rho_b (2K_{\text{sgs}})^{1/2} S_{ij}^* - 4C_2 \rho_b K \frac{U_{i,k} U_{j,k}}{|\nabla \mathbf{U}|^2} - \frac{2}{3} (1 - C_2) \rho_b K \delta_{ij}.$$
 (5.17)

where  $|\nabla \mathbf{U}| := (2U_{i,k}U_{i,k})^{1/2}$  is the norm of the resolved velocity derivative,

$$S_{ij}^* = S_{ij} - \frac{1}{3}\delta_{ij}d = \frac{1}{2}(U_{i,j} + U_{j,i}) - \frac{1}{3}\delta_{ij}U_{k,k}$$
 (5.18)

is the trace-free rate-of strain, and  $\Delta = (\Delta x \ \Delta y \ \Delta z)^{1/3}$  is the grid scale in comoving coordinates. The production and dissipation terms in equation (5.16) are defined as follows:

$$\Sigma = \frac{1}{a}\tau_{ij}S_{ij}, (5.19)$$

$$\varepsilon = \frac{C_{\varepsilon} K^{3/2}}{a\Delta},\tag{5.20}$$

and  $\kappa_{\rm sgs} = C_\kappa \Delta K^{1/2}$  is the SGS diffusivity. Here we assume that the Reynolds number of turbulence is high such that the damping of turbulent eddies by the microscopic viscosity of the fluid occurs entirely on the subgrid scales. Because of the numerical viscosity of PPM, however, part of the numerically resolved kinetic energy will be dissipated directly into internal energy.

# Gravity

In NYX we always compute gravity by solving a Poisson equation on the mesh hierarchy. To make sure this option is chosen correctly, we must always set

## $USE\_GRAV = TRUE$

in the GNUmakefile and

castro.do\_grav = 1
gravity.gravity\_type = PoissonGrav

in the inputs file.

To define the gravitational vector we set

$$\mathbf{g}(\mathbf{x},t) = -\nabla\phi\tag{6.1}$$

where

$$\Delta \phi = \frac{4\pi G}{a} (\rho - \overline{\rho}) \tag{6.2}$$

where  $\overline{\rho}$  is the average of  $\rho$  over the entire domain if we assume triply periodic boundary conditions, and a(t) is the scale of the universe as a function of time.

# Dark Matter Particles

For the moment, assume that we are running in comoving coordinates, with dark matter particles only (no hydro) and that the particles all exist at level 0. These assumptions are encapsulated in the following lines in the inputs file:

```
nyx.use\_comoving = t
nyx.do_dm_particles = 1
amr.max_level = 0
nyx.do_hydro = 0
nyx.do_react = 0
nyx.do_grav = 1
```

#### **Equations** 7.1

If we define  $\mathbf{x}_i$  and  $\mathbf{u}_i$  as the location and velocity of particle i, respectively, then we wish to solve

$$\frac{d\mathbf{x}_i}{dt} = \frac{1}{a}\mathbf{u}_i \tag{7.1}$$

$$\frac{d\mathbf{x}_i}{dt} = \frac{1}{a}\mathbf{u}_i \tag{7.1}$$

$$\frac{d(a\mathbf{u}_i)}{dt} = \mathbf{g}_i \tag{7.2}$$

where  $\mathbf{g}_i$  is the gravitational force evaluated at the location of particle i, i.e.,  $\mathbf{g}_i = \mathbf{g}(\mathbf{x}_i, t)$ .

#### 7.2 Initializing the Particles

There are several different ways in which particles can currently be initialized:

#### 7.2.1 Read from an ASCII file

To enable this option, set

```
nyx.particle_init_type = AsciiFile
nyx.ascii_particle_file = particle_file
```

Here *particle\_file* is the user-specified name of the file. The first line in this file is assumed to contain the number of particles. Each line after that contains

x y z mass xdot ydot zdot

Note that the variable that we call the particle velocity,  $\mathbf{u} = a\dot{\mathbf{x}}$ , so we must multiply  $\dot{\mathbf{x}}$ , by a when we initialize the particles.

## 7.2.2 Read from a binary file

To enable this option, set

```
nyx.particle_init_type = BinaryFile
nyx.binary_particle_file = particle_file
```

As with the ASCII read, the first line in this file is assumed to contain the number of particles. Each line after that contains

x y z mass xdot ydot zdot

Note that the variable that we call the particle velocity,  $\mathbf{u} = a\dot{\mathbf{x}}$ , so we must multiply  $\dot{\mathbf{x}}$ , by a when we initialize the particles.

## 7.2.3 Read from a binary "meta" file

This option allows you to read particles from a series of files rather than just a single file. To enable this option, set

```
nyx.particle_init_type = BinaryMetaFile
nyx.binary_particle_file = particle_file
```

In this case the *particle\_file* you specify is an ASCII file specifying a list of file names with full paths. Each of the files in this list is assumed to be binary and is read sequentially (individual files are read in parallel) in the order listed.

## 7.2.4 Reading SPH particles

For some applications it is useful to initialize the grid data with SPH-type particles. To enable this option, you must set

```
nyx.do_santa_barbara = 1
nyx.init_with_sph_particles = 1
```

The SPH-type particles can then be read in by setting

```
nyx.sph_particle_file = sph_particle_file
```

where  $sph\_particle\_file$  is the user-specified name of the file containing the SPH particles. The type of  $sph\_particle\_file$  must be the same (Ascii, Binary or BinaryMeta) as the dark matter particle file as specified by

```
nyx.particle_init_type =
```

The SPH particles will be discarded by the code once the grid data has been initialized.

#### 7.2.5 Random placement

To enable this option, set

```
nyx.nyx.particle_init_type = Random
```

There are then a number of parameters to set, for example:

```
nyx.particle_initrandom_count = 100000
nyx.particle_initrandom_mass = 1
nyx.particle_initrandom_iseed = 15
```

#### 7.2.6 Cosmological

Using cosmological initial conditions is a three step process:

- 1. Generating a transfer function (e.g. with camb)
- 2. Generating an initial displacement field (with nyx-ic)
- 3. Starting nyx

In the following we will look at each step a bit closer.

#### 7.2.6.1 Generating a transfer function

The transfer function is used in nyx-ic to generate the power spectrum. The usual way is to use camb¹ to calculate it for the desired universe. A sample camb.ini is provided with nyx-ic. The important options are:

- transfer\_redshift(1) = 50
- $transfer_matterpower(1) = tf$

which determine the initial time for the simulation. You should make sure that you catch all necessary wave numbers for the considered box length and resolution.

From the camb output you have to note values for sigma\_8 for a redshift of zero and the initial redshift. We need this to compute the right normalization.

<sup>&</sup>lt;sup>1</sup>See http://camb.info/

#### 7.2.6.2 Setting up the initial displacements

We calculate the initial displacements with a stand-alone program called nyx-ic. This takes a transfer function and some cosmological parameters as an argument and outputs an "init" directory which basically contains initial displacements for every grid point in a BoxLib MultiFAB. Furthermore the mf contains a fourth field containing the density contrast as initial condition for the baryonic matter.

nyx-ic is started with an "inputs" file similar to the one from Nyx. A sample one is provided. The options are

```
#Omega_{Matter}
cosmo.omegam = 0.272
#Omega_{Lambda}
cosmo.omegax = 0.728
#equation of state paramater omega_{effective}
cosmo.weff = -0.980
#Omega_{baryon}*Hubble^2
cosmo.ombh2 = 0.0226
#Hubble/100km/s
cosmo.hubble = 0.704
#scalar spectral index
cosmo.enn = 0.963
# initial z
cosmo.z_init = 50
#sidelength of the box (in Mpc)
cosmo.boxside = 90.14
#seed of the rng
cosmo.isd = 100
#resolution of the box
cosmo.gridpoints = 256
#the output file name
cosmo.initDirName = init
#choose the source of the transferfunction
cosmo.transferfunction = CAMB
#some tabulated transferfunction generated with camb (compare camb-ini-file)
cosmo.tabulatedTk = tf
# sigma8 for the input tf at z=0 and initial z (to calc the growthfactor)
cosmo.init\_sigma8_0 = 0.7891368
cosmo.init_sigma8_init = 2.0463364E-02
```

The code solves the equation

$$P(k,a) = 2\pi^2 \delta_H^2 \frac{k^n}{H_0^{n+3}} T^2(k) \left(\frac{D(a)}{D(a=1)}\right)^2$$
(7.3)

to calculate P and from that gaussian distributed density perturbations  $\delta$  following that spectrum. Particle displacements are then calculated as Zel'dovich displacements.

Non-gaussian effects as well as neutrino contributions are planned for the future.

#### 7.2.6.3 Using Nyx with cosmological initial conditions

- nyx.nyx.particle\_init\_type = Cosmological set the *right* init type
- cosmo.initDirName = init set the name of the displacements directory (boxlib format)
- cosmo.particle\_mass = 0.19178304E+10 sets the mass  $[M_{\odot}]$  of each particle
- cosmo.omegam = 0.272 set  $\Omega_{Matter}$
- cosmo.omegax = 0.728 set  $\Omega_{\Lambda}$
- cosmo.hubble = 0.704 set the reduced hubble constant h

We will generate a particle of mass particle\_mass in every grid cell displaced from the center by the value found in the **initDirName** for that cell. Velocities are calculated in the Zel'dovich approximation by

$$\vec{v} = \Delta \vec{x} \times 100 \text{km/s} \times a \sqrt{\Omega_M / a^3 + \Omega_\Lambda} \times L_{\text{box}}$$
 (7.4)

where  $\Delta \vec{x}$  is the displacement of the particle.

## 7.3 Time Stepping

There are currently two different ways in which particles can be moved:

#### **7.3.1** Random

To enable this option, set

#### nyx.particle\_move\_type = Random

Update the particle positions at the end of each coarse time step using a random number between 0 and 1 multiplied by 0.25 dx.

#### 7.3.2 Motion by Self-Gravity

To enable this option, set

nyx.particle\_move\_type = Gravitational

#### 7.3.2.1 Move-Kick-Drift Algorithm

In each time step:

- Solve for  $\mathbf{g}^n$  (only if multilevel, otherwise use  $\mathbf{g}^{n+1}$  from previous step)
- $\mathbf{u}_{i}^{n+\frac{1}{2}} = \frac{1}{a^{n+1/2}}((a^{n}\mathbf{u}_{i}^{n}) + \frac{\Delta t}{2}\mathbf{g}_{i}^{n})$
- $\bullet \ \mathbf{x}_i^{n+1} = \mathbf{x}_i^n + \frac{\Delta t}{a^{n+1/2}} \mathbf{u}_i^{n+1/2}$
- Solve for  $\mathbf{g}^{n+1}$  using  $\mathbf{x}_i^{n+1}$
- $\bullet \ \mathbf{u}_i^{n+1} = \tfrac{1}{a^{n+1}}((a^{n+\frac{1}{2}}\mathbf{u}_i^{n+\frac{1}{2}}) + \tfrac{\Delta t}{2} \ \mathbf{g}_i^{n+1})$

Note that at the end of the timestep  $\mathbf{x}_i^{n+1}$  is consistent with  $\mathbf{g}^{n+1}$  becase we have not advanced the positions after computing the new-time gravity. This has the benefit that we perform only one gravity solve per timestep (in a single-level calculation with no hydro) because the particles are only moved once.

#### 7.3.2.2 Computing g

We solve for the gravitational vector as follows:

- Assign the mass of the particles onto the grid in the form of density,  $\rho_{DM}$ . The mass of each particle is assumed to be uniformly distributed over a cube of side  $\Delta x$ , centered at what we call the position of the particle. We distribute the mass of each particle to the cells on the grid in proportion to the volume of the intersection of each cell with the particle's cube. We then divide these cell values by  $\Delta x^3$  so that the right hand side of the Poisson solve will be in units of density rather than mass. Note that this is the *comoving* density.
- Solve  $\nabla^2 \phi = \frac{4\pi G}{a} \rho_{DM}$ . We discretize with the standard 7-point Laplacian (5-point in 2D) and use multigrid with Gauss-Seidel red-black relaxation to solve the equation for  $\phi$  at cell centers.
- Compute the normal component of  $\mathbf{g} = -\nabla \phi$  at cell faces by differencing the adjacent values of  $\phi$ , e.g. if  $\mathbf{g} = (g_x, g_y, g_z)$ , then we define  $g_x$  on cell faces with a normal in the x-direction by computing  $g_{x,i-1/2,j,k} = -(\phi_{i,j,k} \phi_{i-1,j,k})/\Delta x$ .
- Interpolate each component of **g** from normal cell faces onto each particle position using linear interpolation in the normal direction.

### 7.4 Output Format

#### 7.4.1 Checkpoint Files

The particle positions and velocities are stored in a binary file in each checkpoint directory. This format is designed for being read by the code at restart rather than for diagnostics.

We note that the value of a is also written in each checkpoint directory, in a separate ASCII file called  $comovinq_{-}a$ , containing only the single value.

#### 7.4.2 Plot Files

If **particles.write\_in\_plotfile** = 1 in the inputs file then the particle positions and velocities will be written in a binary file in each plotfile directory.

In addition, we can also visualize the particle locations as represented on the grid. There are two "derived quantities" which represent the particles. Setting

```
amr.derive_plot_vars = particle_count particle_mass_density amr.plot_vars = NONE
```

in the inputs file will generate plotfiles with only two variables. **particle\_count** represents the number of particles in a grid cell; **particle\_mass\_density** is the density on the grid resulting from the particles.

We note that the value of a is also written in each plotfile directory, in a separate ASCII file called  $comoving_a$ , containing only the single value.

#### 7.4.3 ASCII Particle Files

To generate an ASCII file containing the particle positions and velocities, one needs to restart from a checkpoint file but doesn't need to run any steps. For example, if chk00350 exists, then one can set:

```
amr.restart = chk00350
max_step = 350
particles.particle_output_file = particle_output
```

which would tell the code to restart from chk00350, not to take any further time steps, and to write an ASCII-format file called  $particle\_output$ .

This file has the same format as the ASCII input file:

```
number of particles x y z mass xdot ydot zdot
```

### 7.4.4 Run-time Data Logs

If you set

 $amr.data_log = log_file$ 

in the inputs file, then at run-time the code will write a log file with entries every coarse grid time step, containing

nstep time dt redshift a

### 7.4.5 Run-time Screen Output

There are a number of flags that control the verbosity written to the screen at run-time. These are:

amr.v nyx.v gravity.v mg.v particles.v

These control printing about the state of the calculation (time, value of a, etc) as well as timing information.

## Chapter 8

## Visualization

The BoxLib format in which NYX output is written can be read by amrvis, VisIt, and yt.

#### 8.1 amrvis

We have a homegrown visualization tool called amrvis. We encourage you to build the amrvis3d executable, and to try using it to visualize your data. A very useful feature is View/Dataset, which allows you to actually view the numbers – this can be handy for debugging. You can modify how many levels of data you want to see, whether you want to see the grid boxes or not, what palette you use, etc.

If you like to have amrvis display a certain variable, at a certain scale, when you first bring up each plotfile (you can always change it once the amrvis window is open), you can modify the amrvis.defaults file in your directory to have amrvis default to these settings every time you run it.

#### 8.2 VisIt

**VisIt** is also a great visualization tool, and it directly handles our plotfile format (which it calls Boxlib).

#### See http://visit.llnl.gov

To use the Boxlib3D plugin, select it from File  $\rightarrow$  Open file  $\rightarrow$  Open file as type Boxlib, and then the key is to read the Header file, plt00000/Header, for example, rather than telling to to read plt00000.

#### 8.3 yt

yt also handles BoxLib format and is a great visualization tool for Nyx output.

Here are quick instructions from Matthew Turk:

The directories require that the inputs file be one level up, so that the hierarchy of files looks something like:

```
data/data/inputs
data/plt00001
data/plt00002

To load the data in yt, you would then do:
from yt.mods import *
pf = load("data/plt00001")
```

You can also be in the data/directory and just load plt00001.

See http://yt.enzotools.org to download yt and for more information.

### 8.4 Controlling What's in the PlotFile

```
amr.plot_vars =
```

and

#### $amr.derive\_plot\_vars =$

are used to control which variables are included in the plotfiles. The default for **amr.plot\_vars** is all of the state variables. The default for **amr.derive\_plot\_vars** is none of the derived variables. So if you include neither of these lines then the plotfile will contain all of the state variables and none of the derived variables.

If you want all of the state variables plus entropy and pressure (both derived quantities), for example, then set

```
amr.derive_plot_vars = entropy pressure
```

If you just want density (state variable) and pressure (derived quantity), for example, then set

```
amr.plot_vars = density
```

```
amr.derive\_plot\_vars = pressure
```

Recall that we can also control whether the particles are written into a separate file in the plotfile directory by setting

```
particles.write\_in\_plotfile = 1
```

## Chapter 9

## Software Framework

#### 9.1 Code structure

The code structure in the NYX directory that you have checked out is the following:

- Parallel: the "primary" directory, all in C++/Fortran
  - amrlib: basic routines necessary for AMR
  - bndrylib : basic interface routines
  - BoxLib: the most basic directory, everything depends on classes defined here
  - **Nyx**: this is where all the actual algorithm stuff lives
    - \* Exec : various examples
      - · Test\_90Mpc: run directory for the 90Mpc box problem
    - \* Source : source code
    - \* UsersGuide : you're reading this now!
  - mglib: this is the MultiGrid solver written in C++/Fortran would be used for the Poisson solve if we didn't use the F90 solver
  - MGT solver: this is the interface between the C++ code and the F90 multigrid solver
  - **mk**: makefile stuff for C++/Fortran
  - pAmrvis: contains amrvis, a visualization tool for 2D and 3D plotfiles
  - scripts : compiling stuff for C++/Fortran
  - util: various data analysis utilities
  - volpack: package required to compile and run amrvis in 3d
- **fParallel**: the F90 directory, used here only for the multilevel Poisson solver, EOS, and neworks.
  - boxlib: the most basic directory which defines things for the F90 codes
  - data processing: this contains Fortran routines that read in Nyx plotfiles and can do simply processing, including extracting a line along the x, y, or z-axis, averaging a solution over spherical angles to get the profile as a function of radius, and dumping out a brick of data that can be read by IDL see Mike Zingale if interested

- extern: contains EOS and networks

- mg: F90 multigrid - used for the gravity solver only

- **mk** : makefile stuff for F90

- scripts : compiling stuff for F90

Within Parallel/Nyx are the following files:

• Nyx.cpp: this holds the time advancement algorithm

• Nyx\_setup.cpp: this is where components of the state, boundary conditions, derived quantities, and error estimation quantities are defined for a run

• MacBndry.cpp: this is needed to correctly do the adaptive boundary conditions for the Poisson solver

• main.cpp: initializes the BoxLib and timing stuff properly – don't mess with this

#### 9.2 Variable Names

The following is a list of variables, routines, etc used in NYX. It may not be complete or even entirely accurate; it's mostly intended for my own use.

lo, hi: index extent of the "grid" of data currently being handled by a NYX routine

**domlo, domhi**: index extent of the problem domain. This changes according to refinement level: 0th refinement level will have 0, castro.max\_grid\_size, and nth level will go from 0 to castro.max\_grid\_size\*(multiplying equivalent of sum)castro.ref\_ratio(n).

dx: cell spacing, presumably in cm, since CASTRO uses cgs units

**xlo**: physical location of the lower left-hand corner of the "grid" of data currently being handled by a CASTRO routine

**bc**: array that holds boundary condition of and array. Sometimes it appears of the form bc(:,:) and sometimes bc(:,:,:). The last index of the latter holds the variable index, i.e. density, pressure, species, etc.

## 9.3 Parallel I/O

Both checkpoint files and plotfiles are really directories containing subdirectories: one subdirectory for each level of the AMR hierarchy. The fundamental data structure we read/write to disk is a MultiFab, which is made up of multiple FAB's, one FAB per grid. MultiPab multiFabs may be written to each directory in a checkpoint file. MultiFabs of course are shared across CPUs; a single MultiFab may be shared across thousands of CPUs. Each CPU writes the part of the MultiFab that it owns to disk, but they don't each write to their own distinct file. Instead each MultiFab is

written to a runtime configurable number of files N (N can be set in the inputs file as the parameter amr.checkpoint\_nfiles and amr.plot\_nfiles; the default is 64). That is to say, each MultiFab is written to disk across at most N files, plus a small amount of data that gets written to a header file describing how the file is laid out in those N files.

What happens is N CPUs each opens a unique one of the N files into which the MultiFab is being written, seeks to the end, and writes their data. The other CPUs are waiting at a barrier for those N writing CPUs to finish. This repeats for another N CPUs until all the data in the MultiFab is written to disk. All CPUs then pass some data to CPU 0 which writes a header file describing how the MultiFab is laid out on disk.

We also read MultiFabs from disk in a "chunky" manner opening only N files for reading at a time. The number N, when the MultiFabs were written, does not have to match the number N when the MultiFabs are being read from disk. Nor does the number of CPUs running while reading in the MultiFab need to match the number of CPUs running when the MultiFab was written to disk.

Think of the number N as the number of independent I/O pathways in your underlying parallel filesystem. Of course a "real" parallel filesystem should be able to handle any reasonable value of N. The value -1 forces N to the number of CPUs on which you're running, which means that each CPU writes to a unique file, which can create a very large number of files, which can lead to inode issues.

## Chapter 10

# Verification Test Problems

## 10.1 Cosmology Test Problems

#### 10.1.1 90Mpc Box Problem from Cosmic Data Arxiv

This is a standard test problem with initial particle positions and velocities taken from the Cosmic Data Arxiv (http://t8web.lanl.gov/people/heitmann/arxiv).

The size of the box is 90.14 Mpc on a side.

The initial redshift is z = 50; we end the calculation at redshift = 0; equivalently at a = 1. (Recall that a = 1/(1+z).

# References

- [1] A. Maier, L. Iapichino, W. Schmidt, and J. C. Niemeyer. Adaptively Refined Large Eddy Simulations of a Galaxy Cluster: Turbulence Modeling and the Physics of the Intracluster Medium. *Astrophysical Journal*, 707:40–54, 2009.
- [2] P. Sagaut. Large eddy simulation for incompressible flows: An introduction. Berlin: Springer-Verlag, 2006.
- [3] W. Schmidt and C. Federrath. A fluid-dynamical subgrid scale model for highly compressible astrophysical turbulence. *Astronomy and Astrophysics*, 528:A106+, April 2011.
- [4] W. Schmidt, J. C. Niemeyer, and Hillebrandt. A localised subgrid scale model for fluid dynamical simulations in astrophysics. I. Theory and numerical tests. Astronomy and Astrophysics, 450:265–281, 2006.
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